

## Negligible influence of small enough jitter in the neuron geometry on the extracellular potential

Reviewer 3: supplemental material for the manuscript by Makarova et al., 2011, "Parallel readout of the pathway-specific inputs to laminated brain structures"

### A. Initial definitions

Let

$$\varphi^0(r, t) = -\frac{1}{4\pi\sigma} \sum_{n=1}^N \sum_{c=1}^C \frac{I_{nc}^0(t)}{d_{nc}^0} \quad (1)$$

be the potential created by an aggregate of  $N$  identical, perfectly aligned neurons with  $C$  compartments (see also Eq. (3) in the manuscript). Under the perfect alignment we understand parallel alignment of the neurons over (vertical)  $z$ -axis with no rotation in  $(x, y)$ -plane. Somas of all neurons have the same  $z$  position but can be randomly distributed in  $(x, y)$ -plane. In Eq. (1)  $d_{nc}^0 = \|r_{nc}^0 - r\|$  is the Euclidean distance between the compartment  $nc$  and the recording point,  $I_{nc}^0(t)$  is the transmembrane current generated by the compartment, and  $\sigma$  is the (constant) conductivity of the extracellular space. We notice that

$$d_{nc}^0 \geq \delta > 0, \quad \forall n, c \quad (2)$$

where  $\delta$  is a small enough number, otherwise the model of the potential (1) is not valid. Under the assumption of the same synaptic input Eq. (1) can be simplified

$$\varphi^0(r, t) = -\frac{1}{4\pi\sigma} \sum_{c=1}^C I_c^0(t) R_c^0 \quad (3)$$

where  $R_c^0 = \sum_{n=1}^N \frac{1}{d_{nc}^0}$  is the cumulative inverse distance to the  $c$ -th compartment of all neurons (this sum converges to a constant depended on the mean intersoma distance).

### B. Jitter in the compartment's diameter

Let us consider the case of variation of the diameter of the compartments among different neurons in the aggregate. The diameter of the  $c$ -th compartment of the  $n$ -th neuron is given by:

$$a_{nc} = a_{nc}^0 + \epsilon \xi_{nc} \quad (4)$$

where  $a_{nc}^0$  is the diameter of the original (unperturbed) compartment,  $0 < \epsilon$  is the smallness parameter and  $\xi_{nc}$  is a random zero mean variable describing the perturbation.

The compartmental transmembrane current is proportional to the compartment's diameter:  $I_{nc} \propto a_{nc}$ . Thus the variation (4) induces variation in the compartmental currents of the order of magnitude  $\epsilon$  (to obtained exact values, we should also take into account the readjustment of the currents due to their intracellular interaction). Using the Taylor expansion we can write:

$$I_{nc}(t) = I_{nc}^0(t) + \epsilon J_{nc}(t) + O(\epsilon^2) \quad (5)$$

where  $J_{nc}(t)$  are the transmembrane currents induced by the variation of the diameter of the compartments. Then the new currents given by (5) lead to changes in the potential:

$$\varphi(r, t) = \varphi^0(r, t) + \epsilon \psi(r, t) + O(\epsilon^2) \quad (6)$$

where

$$\psi(r, t) = -\frac{1}{4\pi\sigma} \sum_{n=1}^N \sum_{c=1}^C \frac{J_{nc}(t)}{d_{nc}^0} \quad (7)$$

is the first order correction to the potential. In Eq. (6) we notice that the correction of the potential is of the order  $\epsilon$ . Besides, since  $J_{nc}$  is a random variable of its indexes the sum (7) will be significantly smaller in magnitude than the sum (1). To demonstrate this statement we introduce new random variables:

$$\eta_c = \sum_{n=1}^N \frac{J_{nc}}{d_{nc}^0} \quad (8)$$

Since  $N \gg 1$  ( $N = 16966$  in simulations), the central limit theorem is applied and we have

$$\eta_c \sim \mathcal{N}(\mu_\eta, \lambda_\eta) \quad (9)$$

One can easily show that the mean value of  $\eta_c$  is zero ( $\mu_\eta = 0$ ), whereas the standard deviation is given by

$$\lambda_\eta = \lambda_j \sqrt{\sum_{n=1}^N \frac{1}{(d_{nc}^0)^2}} \quad (10)$$

where  $\lambda_j$  is the standard deviation of the random correcting transmembrane currents  $J_{nc}$ . Assuming that  $\lambda_j$  does not depend on the compartment's number (which is true at least at the first order of magnitude) we then obtain the final expression for the correcting potential:

$$\psi \sim \mathcal{N}(0, \lambda_\psi), \quad \lambda_\psi = \lambda_j \sqrt{C \sum_{n=1}^N \frac{1}{(d_{nc}^0)^2}} \quad (10)$$

Thus small enough jitter in the compartment' diameters makes no significant change in the extracellular potential.